

# Correlation of the Peierls Stresses with Crystal Structures Based on a Generalized Peierls-Nabarro Model

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The Peierls stress  $\tau_p$ , which is one of the most fundamental quantities describing the plastic behaviors of crystalline solids, is the minimum external shear stress necessary for a dislocation to overcome the potential barrier arising from the lattice periodicity without any assistance of thermal or quantum lattice vibration. Most theoretical studies on the Peierls stress of dislocations in crystals is based on the Peierls-Nabarro (PN) dislocation model, where only two lattice planes facing the glide plane are treated as discrete lattice and the two half infinite crystals above and below the planes are approximated as elastic continuum. The formula, obtained by Peierls and Nabarro[1,2] and revised by Huntington[3] assuming the sinusoidal shear interaction between atomic rows in the two planes, is  $\tau_p/G=C\exp(Ah/b)$  with  $A=\pi/(1-\nu)$  and  $C=1/(1-\nu)$ , where  $G$  is the shear modulus,  $\nu$  the Poisson ratio,  $h$  the glide plane spacing and  $b$  the magnitude of the Burgers vector. This formula connects the parameters describing the crystal structure to the value of  $\tau_p$ . Ohsawa et al.[4] computed self-consistently the equilibrium configuration of the PN dislocation assuming the longer-range inter-atomic-row potentials instead of the sinusoidal one, and showed that the above formula is satisfied approximately but with different  $A$  and  $C$  values depending on the shape of the interaction potential.

The PN model and its formulation have so far been described in literatures mainly for the edge dislocation in the lattice with the periodicity  $b$  in the direction of dislocation motion. However, in many real crystals, it is known that non-edge dislocation controls the deformation and the periodicity is not same as  $b$ . Hence, in order to compare experimentally estimated  $\tau_p$ 's with the PN models, it is necessary to generalize the PN formula for dislocations with a variety of characters. In the generalization,  $b$  in the original formula should be replaced by lattice periodicity  $\delta$  in the direction of the dislocation motion. The factors  $A$  and  $C$  were obtained to be  $A=\pi/(1-\nu)$  and  $C=1/(1-\nu)$  for edge dislocation and  $A=\pi$  and  $C=1$  for screw dislocation.  $A$  and  $C$  for mixed dislocations should be given as the conventional weighted averages of the factors for screw and edge dislocations, i.e.  $A=\pi[\sin\theta/(1-\nu)+\cos\theta]$  and  $C=A/\pi$  where  $\theta$  is the angle between the Burgers vector and the direction of the dislocation. Using these factors, the generalized formula of the PN relation is expressed as  $\tau_p/(GC) = \exp(Ah/\delta)$ .

Because the value of  $\tau_p$  corresponds to the critical resolved shear stress  $\tau_c$  at the absolute zero, we could obtain it by extrapolating the temperature  $T$  dependence of  $\tau_c$  directly to the absolute zero in the case of metallic and some ionic crystals for which  $\tau_c$  vs.  $T$  relations are available down to low enough temperatures. However, for materials like covalent crystals and intermetallic compounds,  $\tau_c$ 's have not been measured at sufficiently low  $T$  because of the brittleness at low temperatures. We have estimated  $\tau_p$ 's for such brittle materials by a method based on the line tension model of the kink-pair

formation enthalpy. By this method  $\tau_p$  is given as  $(mk_B T_0)^2 / (d^3 b \kappa)$ , where  $m$  is a constant around 30,  $k_B$  the Boltzmann's constant,  $T_0$  the temperature at which the  $T$ -dependence of  $\tau_c$  vanishes,  $d$  the period of the potential called the Peierls potential which a dislocation feels, and  $\kappa$  the line tension of the dislocation. The value of  $\kappa$  for each crystal can be numerically computed by use of the elastic constants. Despite some ambiguity of  $m$  and  $T_0$ ,  $\tau_p$ 's could be estimated with an accuracy.

$\tau_p$ 's of dislocations of 66 slip systems in 52 crystals have been estimated experimentally either by direct extrapolation of the  $\tau_c$  vs.  $T$  curve to 0 K or by the line-tension approach.  $\tau_p/G$  values are distributed over four orders of magnitude,  $10^{-5}$  to  $10^{-1}$ , but those for a group of crystals with the same crystal structure are within an order of magnitude. In order to see the correlation between  $\tau_p$ 's and the crystal structure, we plot  $\tau_p/(GC)$  against  $Ah/\delta$  in FIG. 1, based on the generalized PN model mentioned above. Generally  $\tau_p$ 's are lower than the theoretical relations of Huntington (H) and Ohsawa et al. ( $O_i$ ) This can be explained partially by the fact that only an incomplete relaxation of dislocation core is taken into consideration in the model. However, the large deviations of about two orders of magnitude at  $Ah/\delta \sim 5$  cannot be attributed solely to it, and some other factors should be involved. In the plot, there is a tendency that  $\tau_p$ 's become higher as the covalent bonding character increases, reflecting the general tendency of the normalized theoretical shear strength of crystals [5].

This work was financially supported by Grant-in-Aid for Scientific Research on Priority Areas "Nano Materials Science for Atomic Scale Modification 474" from Ministry of Education, Culture, Sports, Science and Technology (MEXT) of Japan.

## References

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FIG. 1. Generalized PN plot of the estimated  $\tau_p$ . Lines denoted as H and  $O_i$  are the theoretical relations of Huntington [3] and Osawa et al.[4], respectively.

